program *DIFABS* (Walker & Stuart, 1983) did not operate well with the data. The versions we used do not correctly handle indices over 50. So absorption corrections were performed using a modified version of *DATAP* (Coppens, Lieserowitz & Rabinovich, 1965) and the structure was refined using *SHELXL*93 (Sheldrick, 1993).

Data collection: CAD-4 EXPRESS (Enraf-Nonius, 1993). Cell refinement: CAD-4 EXPRESS. Data reduction: CADAK (Savariault, 1991a). Program(s) used to solve structure: SHELXS86 (Sheldrick, 1985). Program(s) used to refine structure: SHELXL93 (Sheldrick, 1993). Molecular graphics: OR-TEPII (Johnson, 1976). Software used to prepare material for publication: TABPUB (Savariault, 1991b).

Lists of structure factors, anisotropic displacement parameters and complete geometry have been deposited with the IUCr (Reference: DU1116). Copies may be obtained through The Managing Editor, International Union of Crystallography, 5 Abbey Square, Chester CH1 2HU, England.

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NbAs₂

WOLFGANG BENSCH AND WOLFRAM HEID

Institut für Anorganische Chemie, Universität Frankfurt, Marie-Curie Strasse 11, D-60439 Frankfurt, Germany

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Abstract

The structure of niobium diarsenide, NbAs₂, has been determined by single-crystal X-ray diffraction; it was found to crystallize with the OsGe₂ structure type.

Comment

The lattice parameters as well as the possible space groups of NbAs₂ were first reported in 1964 (Saini, Calvert & Taylor, 1964). A second study using the Guinier technique was performed in 1964 (Furuseth & Kjekshus, 1964a,b). In further investigations (Furuseth & Kjekshus, 1965), the crystal structure was solved in space group C2 using integrated intensities from Weissenberg photographs. Problems occurring during the refinement procedure were attributed to absorption effects. It was suggested by Pearson (1967, 1972) that OsGe₂ may be of the NbAs₂ type. A reinvestigation of a number of structures crystallizing with either the $OsGe_2$ or NbAs₂ structures showed that C2/m gives the best refinement results (Calvert, 1992). These structures may be classified in C2/m with OsGe₂ as the prototype structure (Calvert, 1992). Very recently, single crystals of NbSb₂ were reported to crystallize in space group C2/m with the OsGe₂ structure type (Rehr & Kauzlarich, 1994). To confirm that NbAs₂ also crystallizes with the OsGe₂ structure type, we refined our single-crystal data in both C2/m and C2. The final R(F) and wR(F)values were slightly lower for the C2/m refinement (3.03) and 2.88%) compared with those of the C2 refinement (3.13 and 3.12%). The positional parameters differ from those reported by Furuseth & Kjekshus (1964b, 1965). Consequently, the bond distances (Table 2) are also different from the values given in the previous contributions (Furuseth & Kjekshus, 1964b, 1965).

The Nb atom is in a triangular prism of six As atoms. Two other As atoms and one Nb atom lie outside the rectangular faces. Outside each of the triangular faces two, more distant, Nb atoms are located. Each As1 atom has a distorted square-pyramidal environment of five Nb atoms. One close As1 atom is located across the base of the pyramid. Each As2 atom is coordinated to three other As2 and three Nb atoms. The arrangement of the

NbAs₂

As2 atoms leads to alternating short and long As2— As2 distances. The short As2—As2 distances can be regarded as covalent bonds.

Fig. 1. The crystal structure of NbAs₂. Ellipsoids are drawn at the 70% probability level.

Experimental

During our investigation of the quarternary system Tl-Nb-Se-As we obtained NbAs₂ as black needles. The elements in the ratio 0.8:6:7:1 (Tl:Nb:Se:As) were placed in a quartz ampoule and sealed under vacuum. The mixture was heated to 1273 K (100 K h⁻¹) for six weeks and cooled to room temperature by quenching the ampoule in water. Crystals of NbAs₂ were obtained in addition to different niobium selenides.

Crystal data

NbAs ₂	Mo $K\alpha$ radiation		
$M_r = 242.8$	$\lambda = 0.7107 \text{ Å}$		
Monoclinic	Cell parameters from 26		
C2/m	reflections		
a = 9.368(2) Å	$\theta = 10.3 - 18.1^{\circ}$		
b = 3.396(1) Å	$\mu = 35.479 \text{ mm}^{-1}$		
c = 7.799 (3) Å	T = 295 K		
$\beta = 119.42 (1)^{\circ}$	Needle		
$V = 216.1(1) \text{ Å}^3$	$0.14 \times 0.02 \times 0.015$ mm		
Z = 4	Black		
$D_x = 7.460 \text{ Mg m}^{-3}$			

Data collection

Stoe AEDII diffractometer $\omega/2\theta$ scans

294 observed reflections $[F > 4\sigma(F)]$

Absorption correction:
refined from
$$\Delta F$$
 (XEMP
in SHELXTL-Plus;
Sheldrick, 1991)
 $T_{min} = 0.530, T_{max} =$
0.623

664 measured reflections 364 independent reflections

Refinement

Refinement on F R = 0.0303 wR = 0.0363 S = 0.9536 292 reflections 20 parameters w = $1/[\sigma^2(F) + 0.001F^2]$ $(\Delta/\sigma)_{max} \le 0.001$ $R_{int} = 0.0235$ $\theta_{max} = 30^{\circ}$ $h = -12 \rightarrow 0$ $k = -4 \rightarrow 4$ $l = -9 \rightarrow 10$ 4 standard reflections frequency: 120 min intensity decay: <0.5%

 $\begin{aligned} &\Delta \rho_{\text{max}} = 2.84 \text{ e } \text{\AA}^{-3} \\ &\Delta \rho_{\text{min}} = -3.25 \text{ e } \text{\AA}^{-3} \\ &\text{Extinction correction: } F^* = \\ &F(1 + 0.002\chi F^2/\text{sin}\theta)^{-1/4} \\ &\text{Extinction coefficient:} \\ &\chi = 0.0014 (3) \\ &\text{Atomic scattering factors} \\ &\text{from International Tables} \\ &\text{for X-ray Crystallography} \\ &(1974, \text{Vol. IV}) \end{aligned}$

Table 1. Fractional atomic coordinates and equivalent isotropic displacement parameters (Å²)

$$U_{\text{eq}} = (1/3) \sum_i \sum_j U_{ij} a_i^* a_j^* \mathbf{a}_i \cdot \mathbf{a}_j.$$

	x	у	Z	U_{eq}
Nb	0.1574 (1)	0	0.1965 (1)	0.003 (1)
As1	0.4059 (1)	0	0.1084(1)	0.004 (1)
As2	0.3600 (1)	1/2	0.4720 (1)	0.004 (1)

Table 2. Selected distancess (Å)

Nb—As1		2.731 (2)	Nb—Nb		3.034 (1)
Nb—As1	×2	2.705(1)	As1—As1		2.983 (3)
Nb—As2		2.669 (2)	As2—As2		2.439 (2)
Nb—Nb	×2	3.396(1)	As1—As1	×2	3.069 (2)
Nb—As2	×2	2.664 (1)	As2—As2	×2	2.864 (2)
Nb-As1	×2	2.739 (1)			

The *SHELXTL-Plus* package (Sheldrick, 1991) was used for crystallographic computations.

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Lists of structure factors, anisotropic displacement parameters and complete geometry have been deposited with the IUCr (Reference: DU1121). Copies may be obtained through The Managing Editor, International Union of Crystallography, 5 Abbey Square, Chester CH1 2HU, England.

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$NaM^{II}Zr_2F_{11}$ ($M^{II} = Mn$, Fe, Co, Ni, Zn) Fluorides

MOHAMMED HAKAM KETTANI, DANIEL AVIGNANT AND JACQUES METIN

Laboratoire de Chimie des Solides, URA 444, CNRS, Université Blaise Pascal, 63177 Aubière CEDEX, France

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Abstract

The structures of NaMnZr₂F₁₁ (sodium manganese undecafluorodizirconate), NaFeZr₂F₁₁, NaCoZr₂F₁₁, NaNiZr₂F₁₁ and NaZnZr₂F₁₁ have been refined from single-crystal X-ray diffraction data. These compounds are isostructural with AgPdZr₂F₁₁. They are characterized by isolated compressed octahedra, which are slightly tetragonally distorted, surrounding the divalent cations of the first-row transition metals. Attempts at the growth of single crystals of divalent chromium (d^4) and copper (d^9) failed, probably because these cations are unstable as a result of the Jahn–Teller effect.

Comment

Compounds of the general formula NaM^{II}Zr₂F₁₁ can be divided into two very different structure types according to the nature of the divalent cation, M^{2+} . For a divalent cation of the second transition series, such as Pd²⁺ or Ag²⁺, a three-dimensional layered structure with triclinic space group $P\bar{1}$ has been observed (Müller, 1987). In this structure type, the Pd²⁺ ions have regular octahedral coordination, whereas the F-atom environment about Ag²⁺ is pseudotetragonally elongated. It can be inferred that the distortion of the Ag²⁺ octahedra is Jahn-Teller induced. A slightly less distorted variant of this structure type has been observed for $AgPdZr_2F_{11}$, which crystallizes in the monoclinic space group C2/m (Müller, 1987). Here, the Pd²⁺ ion also has regular octahedral coordination. We report herein the structures of a series of fluorides with general formula $NaM^{II}Zr_2F_{11}$ (M^{II} = Mn, Fe, Co, Ni, Zn) which are closely related to that of $AgPdZr_2F_{11}$. In these compounds, the divalent cations are in slightly compressed octahedra. Assuming that the Jahn–Teller effect may operate in Na M^{II} Zr₂F₁₁ compounds since such a strong effect has already been observed for Ag^{2+} in $NaAgZr_2F_{11}$, the lack of representatives for Cr^{2+} (d^4) and Cu^{2+} (d^9), for which a strong Jahn–Teller effect could be expected, may be interpreted as due to such an effect being stronger for divalent cations of the first transition row than for Ag^{2+} , leading to unstable structures. For the Co^{2+} compounds the stabilization of the high-spin state is related to the weak crystal field generally induced by an F-atom environment.

The Zr^{IV} ions are located in corner-linked *trans* chains of $(Zr_2F_{12})^{4-}$ binuclear units comprising pairs of edgesharing pentagonal bipyramids. The Na⁺ and M^{2+} ions are ordered within chains of edge-linked polyhedra, made up of alternating hexagonal bipyramids and octahedra. Both types of chains are joined by sharing corners and edges, resulting in layers as illustrated in Fig. 1. These layers are perpendicular to [201] and are further linked to form a three-dimensional framework by sharing corners. The layers are stacked along [201] in the sequence A, B, C, D, E and are related by a translation in such a way that one layer is displaced relative to the other by $\frac{1}{2}\mathbf{b} - \frac{1}{4}\mathbf{c}$ so that an additional Na–Zr– M^{II} –Zr ordering occurs within the corner-linked *trans* chains of composition (Na $M^{II}Zr_2F_{24})_n^{13n-}$, extended along [201].



Fig. 1. Polyhedral layer showing the linkage between octahedra (M^{II}) , pentagonal bipyramids (Zr) and hexagonal bipyramids (Na).

Experimental

Single crystals of NaMnZr₂F₁₁ were obtained by heating a stoichiometric mixture of the dried binary fluorides at 1073 K for 2 d in a sealed platinum tube followed by slow cooling to room temperature at 4 K h⁻¹. For the other NaM^{II}Zr₂F₁₁ fluorides (M^{II} = Fe, Co, Ni, Zn), single crystals were obtained by heating stoichiometric mixtures of the dried starting fluorides to 1073 K in sealed platinum tubes, as above, but once this temperature had been reached the samples were allowed to cool slowly to 913 K at 4 K h⁻¹. The temperature